

Counterflow-induced inverse energy cascade in three-dimensional superfluid turbulence

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Superfluidity of liquid helium

- Helium-4 displays **superfluidity** below $T_\lambda \approx 2.17$ K
- At $1 \lesssim T < T_\lambda$: **two-fluid phenomenology**
→ coupling between **normal** (viscous) and **superfluid** (inviscid) components

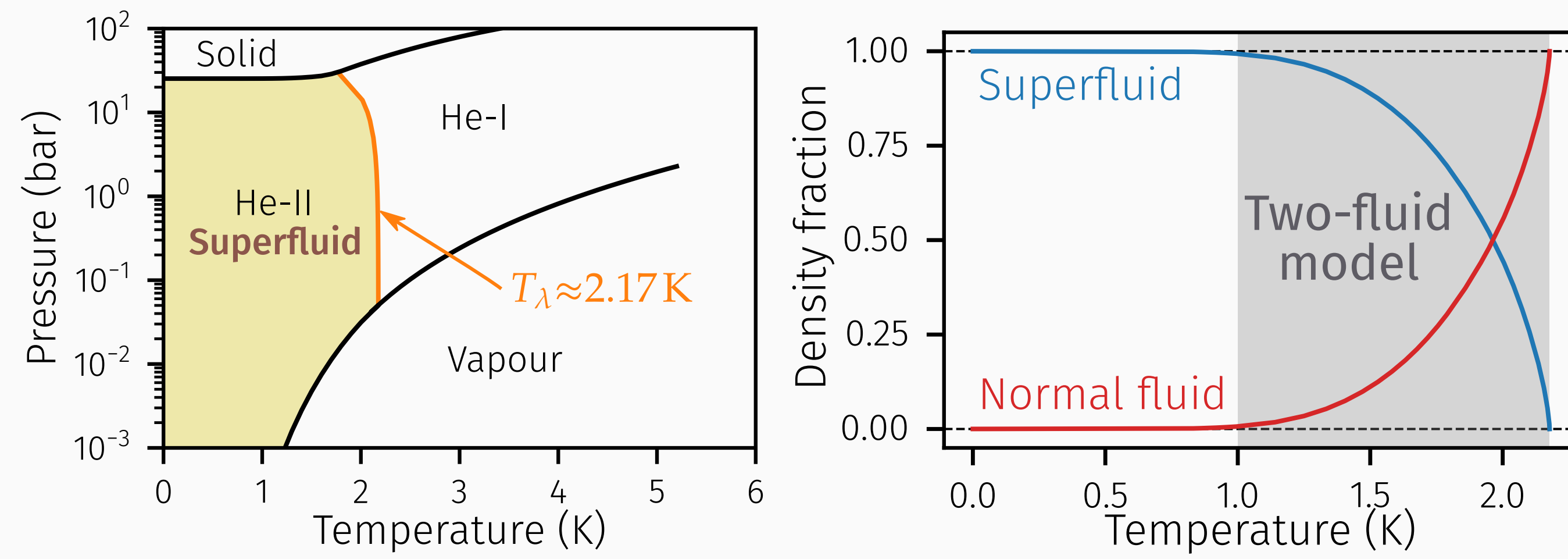


Figure 1. Left: Phase diagram of ^4He . Right: Relative densities of both components.

Counterflow phenomenon

- Produced by a **temperature gradient** in superfluid helium
- Normal** and **superfluid** components flow with **mean relative velocity** \mathbf{U}_{ns}
- Very efficient **heat transport** by normal component

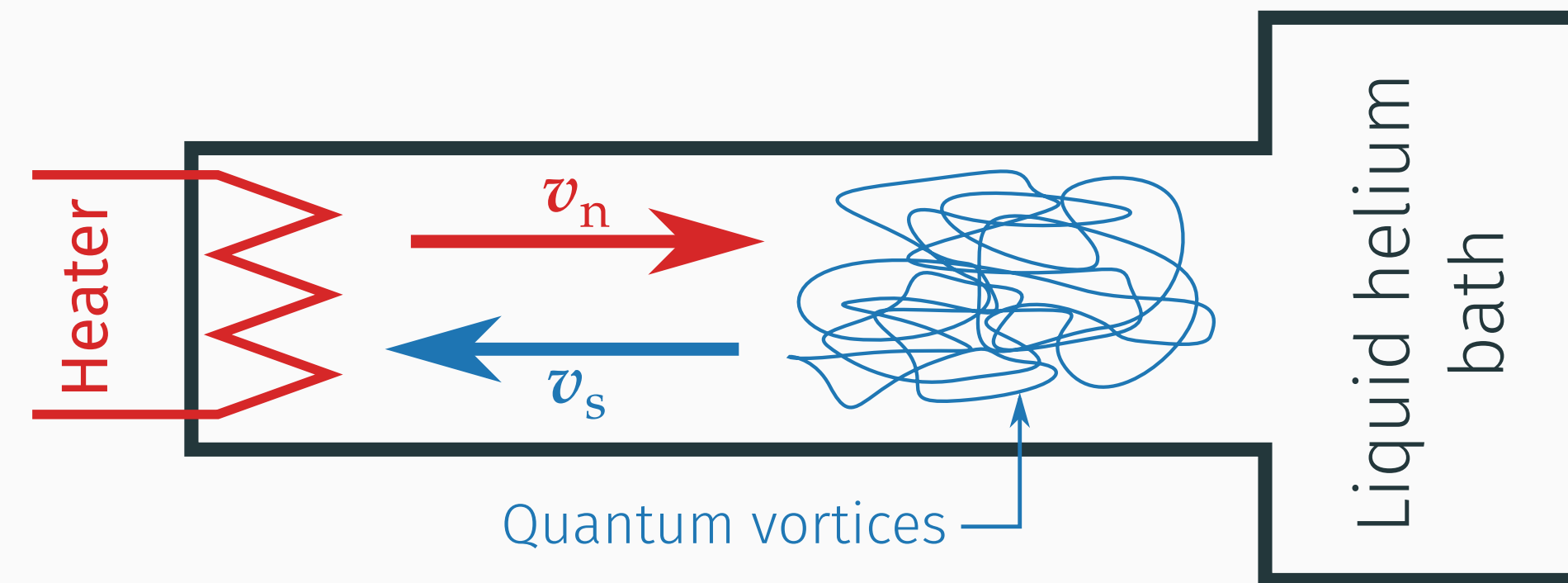


Figure 2. Counterflow in channel of superfluid helium.

Hall-Vinen-Bekarevich-Khalatnikov model

Two-fluid model at macroscopic scales:

$$\begin{aligned} \frac{\partial v_n}{\partial t} + (\mathbf{U}_n + v_n) \cdot \nabla v_n &= -\frac{\nabla p_n}{\rho_n} + \nu_n \nabla^2 v_n - \frac{\mathbf{F}_{\text{ns}}}{\rho_n} + \varphi_n \\ \frac{\partial v_s}{\partial t} + (\mathbf{U}_s + v_s) \cdot \nabla v_s &= -\frac{\nabla p_s}{\rho_s} + \nu_s \nabla^2 v_s + \frac{\mathbf{F}_{\text{ns}}}{\rho_s} + \varphi_s \\ \nabla \cdot v_n &= \nabla \cdot v_s = 0 \end{aligned}$$

- $\mathbf{U}_n, \mathbf{U}_s$ mean velocity of each component ($\Rightarrow \mathbf{U}_{\text{ns}} = \mathbf{U}_n - \mathbf{U}_s$)
- v_n, v_s fluctuating velocity of each component
- $\nu_s(T)$ effective superfluid viscosity → microscopic dissipation
- φ_n, φ_s external forcing $\sim \text{Normal}(0, \sigma_f)$ at wave numbers $|k| \approx k_f$

- $\mathbf{F}_{\text{ns}} = \alpha \rho_s \Omega_0 (v_n - v_s)$ mutual friction force
- $\Omega_0 \sim \sqrt{\langle |\omega_s|^2 \rangle}$ mutual friction frequency
- $\alpha(T)$ temperature-dependent mutual friction parameter

Bidimensionalisation under intense counterflow

- At strong counterflow velocity \mathbf{U}_{ns} :
→ system becomes **quasi-two-dimensional**
→ reminiscent of rotating turbulence, MHD, thin-layer flows, ...

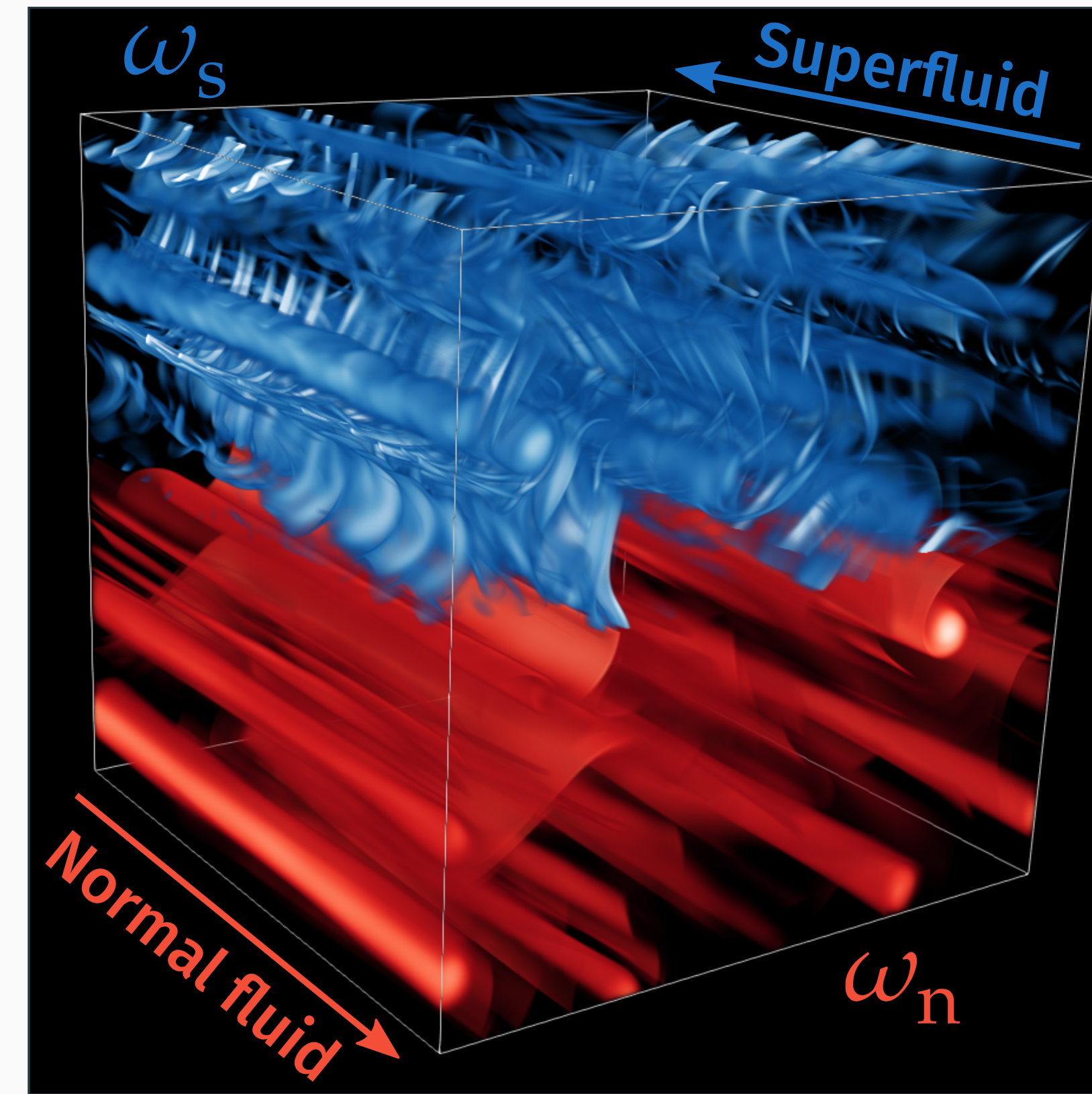


Figure 3. Statistically-steady 3D superfluid turbulence under strong counterflow at $T = 1.9$ K.

Emergence of a split energy cascade

- Starting from initial state with zero fluctuations ($v_n = v_s = 0$):
→ **inverse energy cascade** for $k < k_f$, with $E(k) \sim k^{-5/3}$
→ **direct energy cascade** for $k > k_f$, with $E(k) \sim k^{-3}$
→ scalings compatible with phenomenology of **classical 2D turbulence**

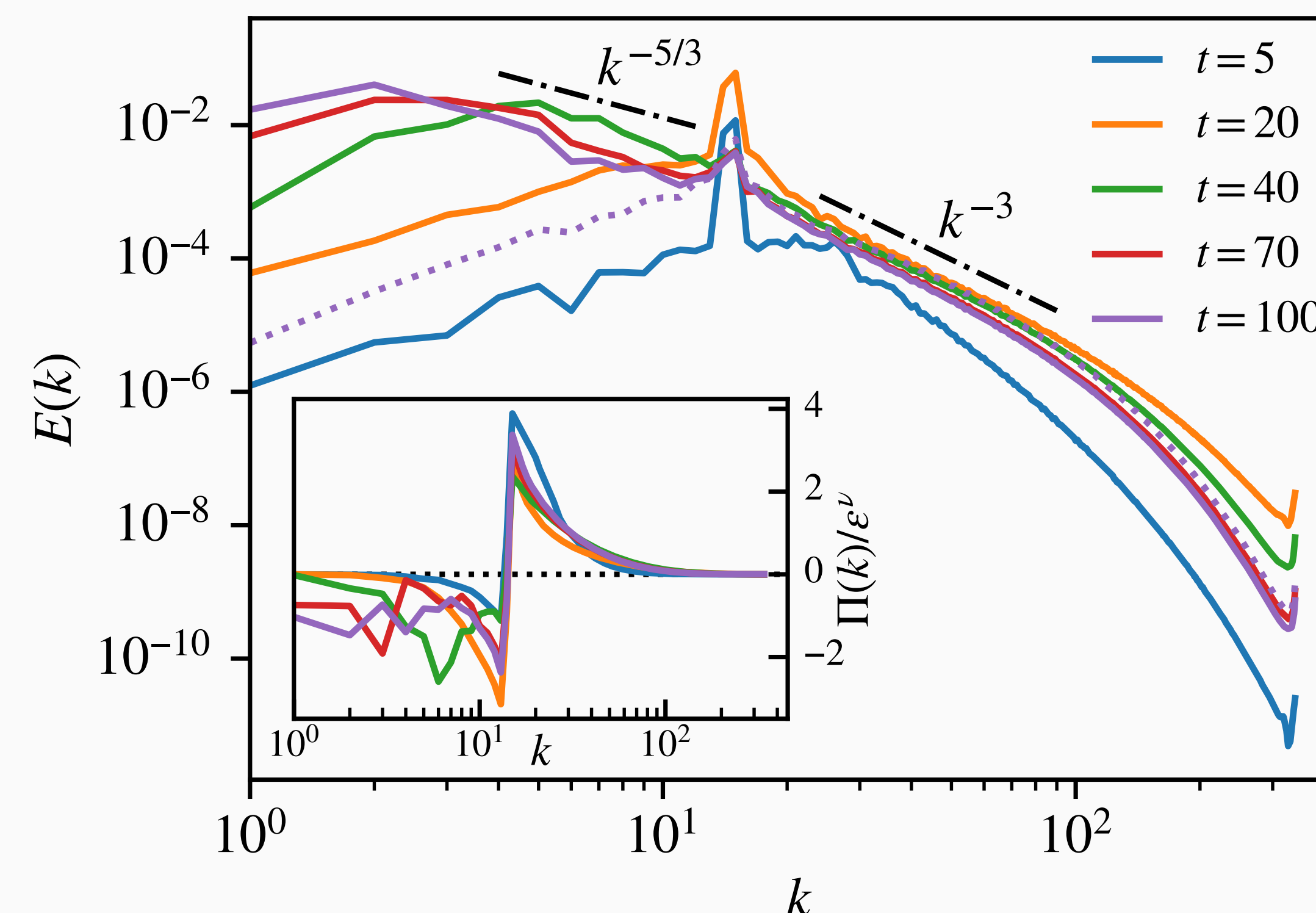


Figure 4. Temporal evolution of kinetic energy spectrum. Times are normalised by the forcing time scale $t_f = (k_f \sigma_f)^{-1/2}$. Inset: normalised energy flux. Simulation performed using $N^3 = 1024^3$ collocation points.

Abrupt transition towards split cascade scenario

Transition happens **abruptly** at critical counterflow velocity U_{ns}^* .

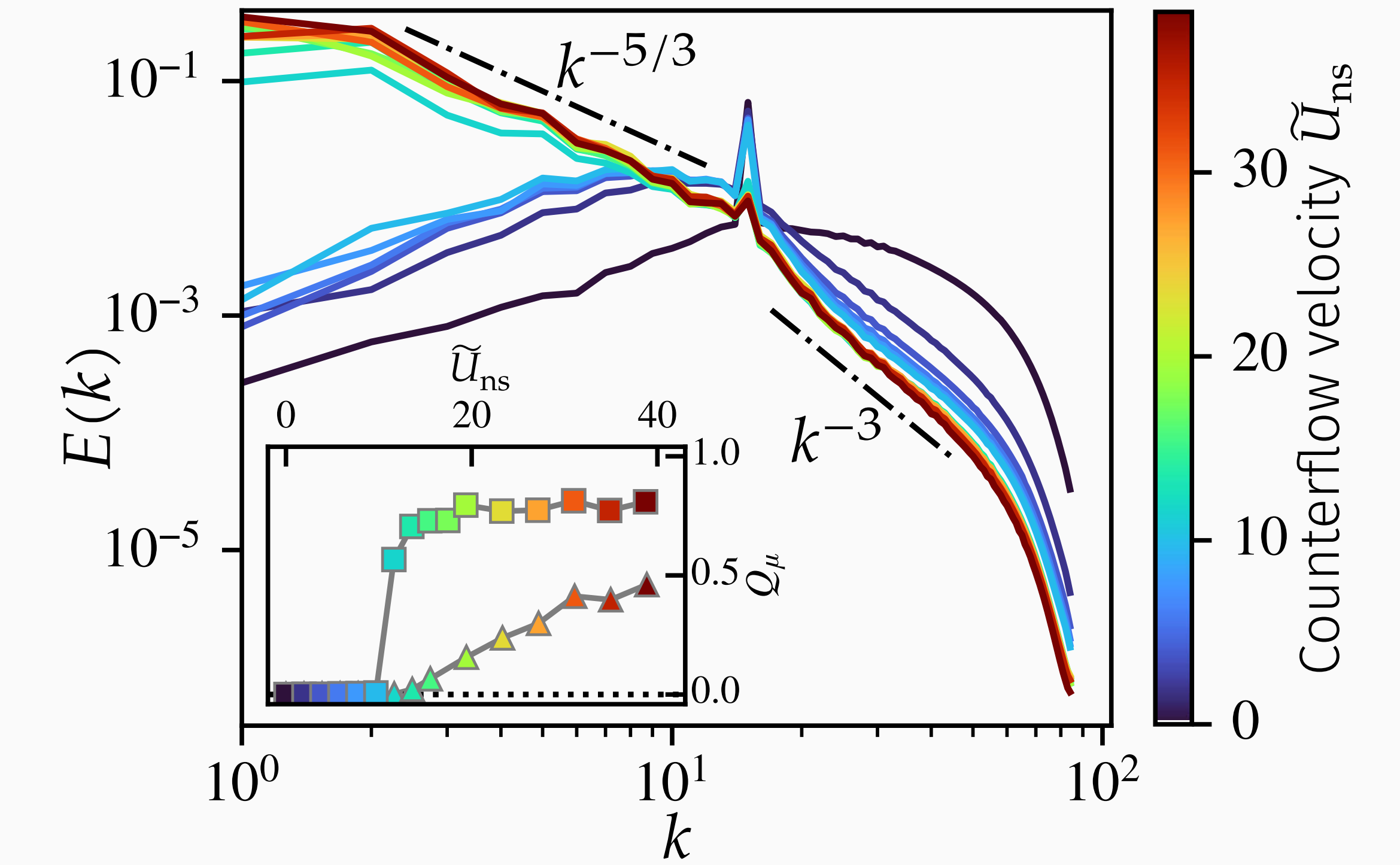


Figure 5. Kinetic energy spectrum for different counterflow velocities. Inset: relative large-scale dissipation using 2D (squares) and 3D (triangles) forcing schemes.

Here, simulations are performed with a **large-scale dissipation term** and with **hyperviscous** small-scale dissipation.

A **2D forcing scheme** is used to obtain a cleaner quasi-2D state at large U_{ns} .

The **critical counterflow velocity** U_{ns}^* can be expressed in terms of the **forcing** and the **mutual friction** parameters.

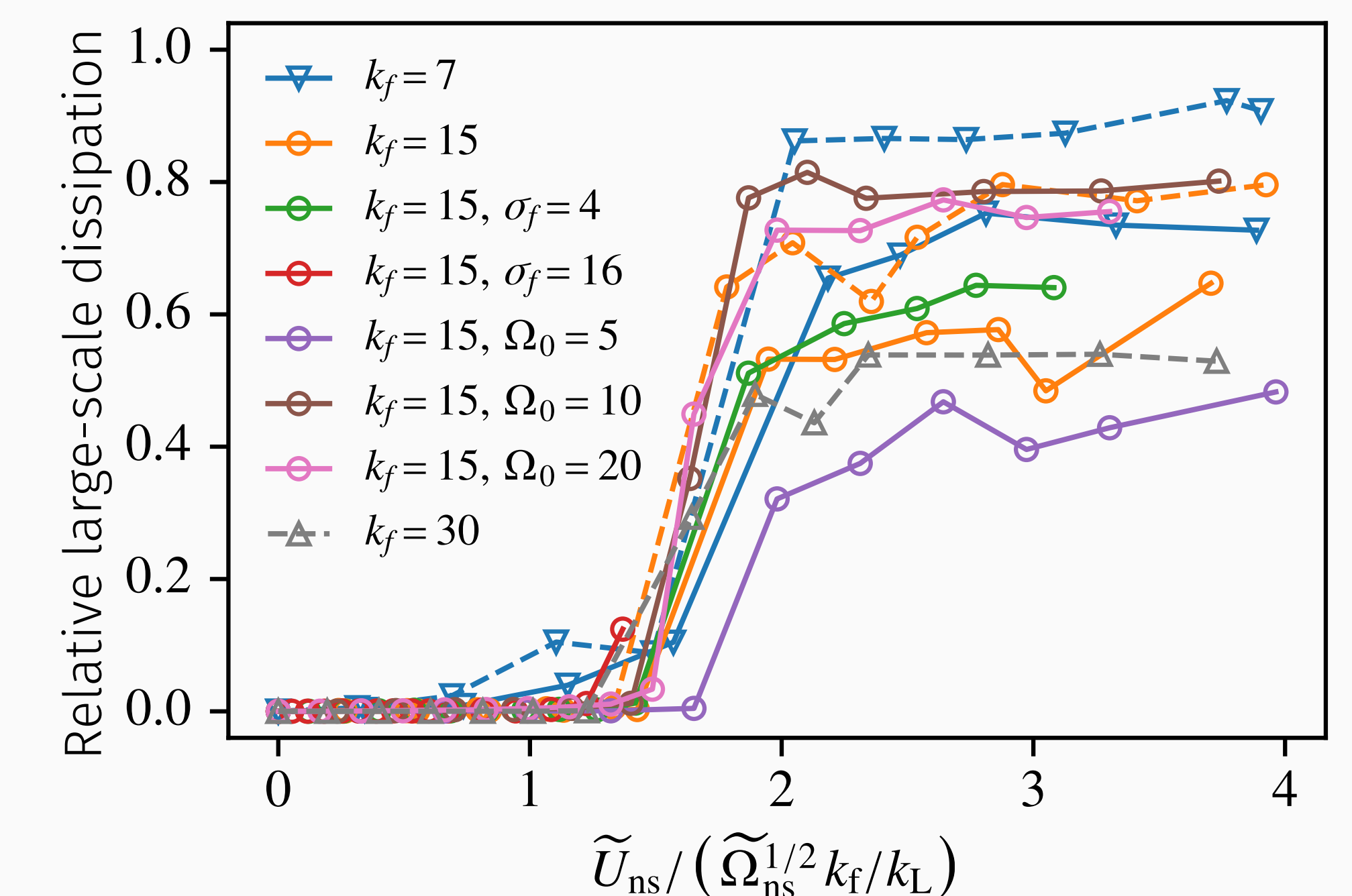


Figure 6. Relative large-scale dissipation Q_L for different forcing and mutual friction parameters.

Perspectives

- Motivate **experimental study** of (quasi-)2D turbulence in superfluid helium.
- The **physical origin** of this transition is not currently understood.
- Characterisation of **temperature effects** and **hysteresis**.
- Application to other **two-fluid systems** (e.g. partially-ionised MHD).